

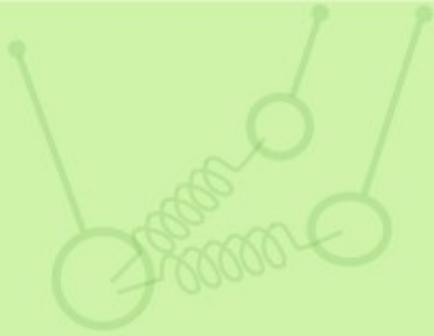


Control of synchronization bistability in oscillatory networks

Irmantas Ratas (irmantas.ff.vu@gmail.com), Kęstutis Pyragas

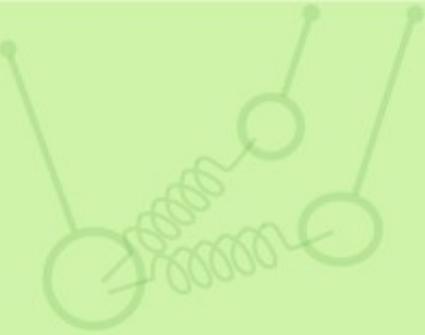
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XXXIV Dynamics Days Europe, Bayreuth, 2014



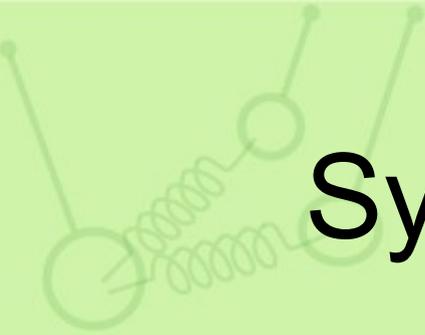
Outline

- Motivation
- Synchronization estimation
- Synchronization bistability
- Algorithms
- Results
- Conclusion



Motivation

- Synchronization – widely observed phenomena
- Pathological synchronization – symptoms of neurological diseases
- Synchronized state – may be not uniquely stable
- Desynchronization methods:
 - I) open loop (e.g. coordinates reset, high frequency stimulation)
 - II) closed loop (e.g. PID, delayed feedback, act-and-wait)



Synchronization estimation

If oscillator have well defined phase, then synchronization is estimated by the ***order parameter***:

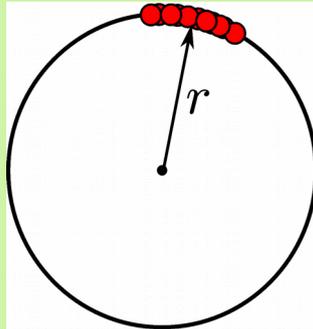
$$r_1 = \frac{1}{N} \sum_{j=1}^N e^{i\theta_j}$$

Synchronization estimation

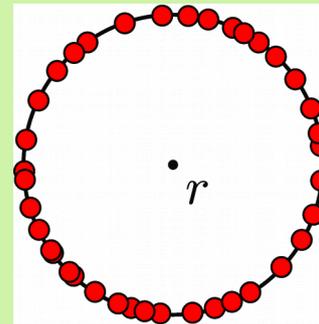
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$|r_1| = 1$
synchronized state



$|r_1| = 0$
desynchronized state

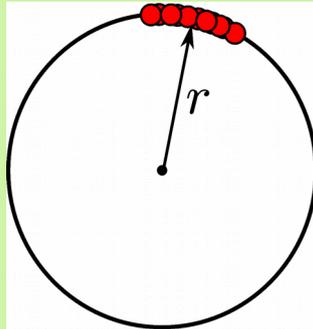


Synchronization estimation

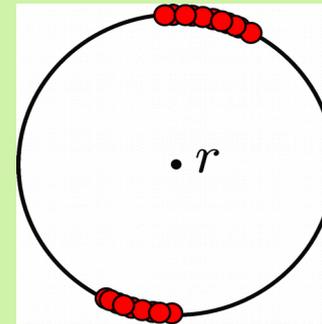
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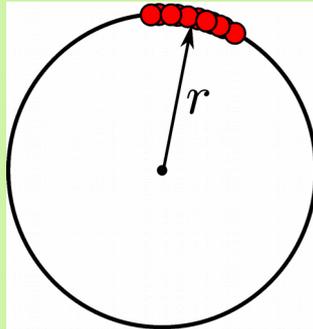
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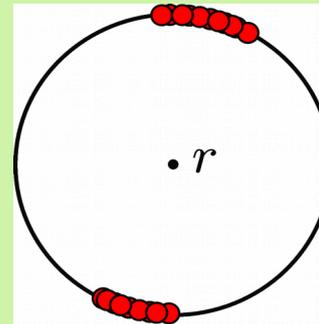
$$r_1 = \frac{1}{N} \sum_{j=1}^N e^{i\theta_j}$$

$$r_n = \frac{1}{N} \sum_{j=1}^N e^{in\theta_j}$$

$|r_1| = 1$
synchronized state



$|r_1| = 0$
desynchronized state



Synchronization estimation

In real life to define separate oscillators phase in coupled network is impossible.

Neuron network models have shape

$$\begin{array}{l} \text{Oscill. Eqs.} \\ \frac{dx_i}{dt} = F(x_i, \mathbf{y}_i) + W(x_1, \dots, x_N), \\ \frac{dy_i}{dt} = \mathbf{G}(x_i, \mathbf{y}_i) \end{array}$$

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Synchronization may be estimated by the **mean potential field variation**:

$$S = \text{Var} \left[\frac{1}{N} \sum_{i=1}^N x_i \right]$$

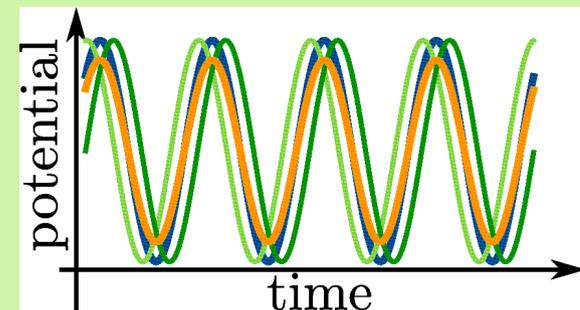
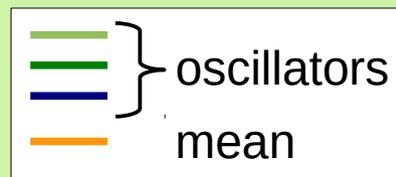
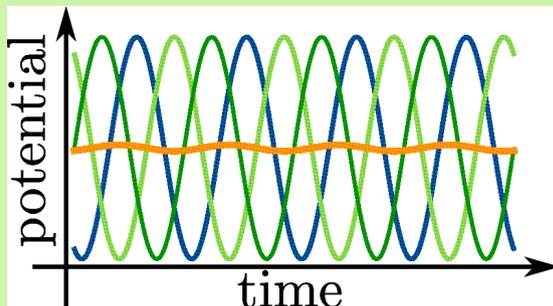
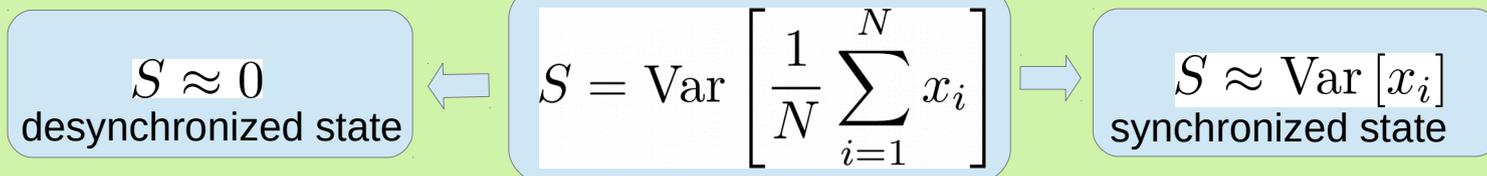
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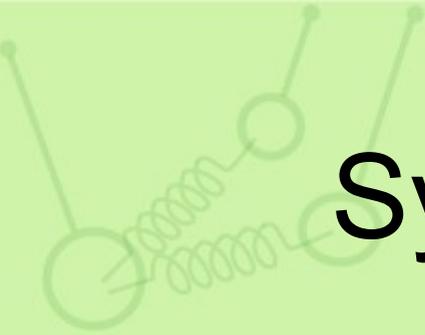
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Synchronization may be estimated by the **mean potential field variation**:





Synchronization bistability

Incoherent and partially synchronized stable states coexist at a particular range of parameters.

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Example:

“Kuramoto Model of Coupled Oscillators with Positive and Negative Coupling Parameters: An Example of Conformist and Contrarian Oscillators” by H. Hong and S. H. Strogatz (PRL, 2011)

$$\dot{\theta}_j = \omega_j + \frac{K_1}{N} \sum_{k=1}^{N_1} \sin(\theta_k - \theta_j) - \frac{K_2}{N} \sum_{k=N_1+1}^N \sin(\theta_k - \theta_j)$$

conformists

contrarians

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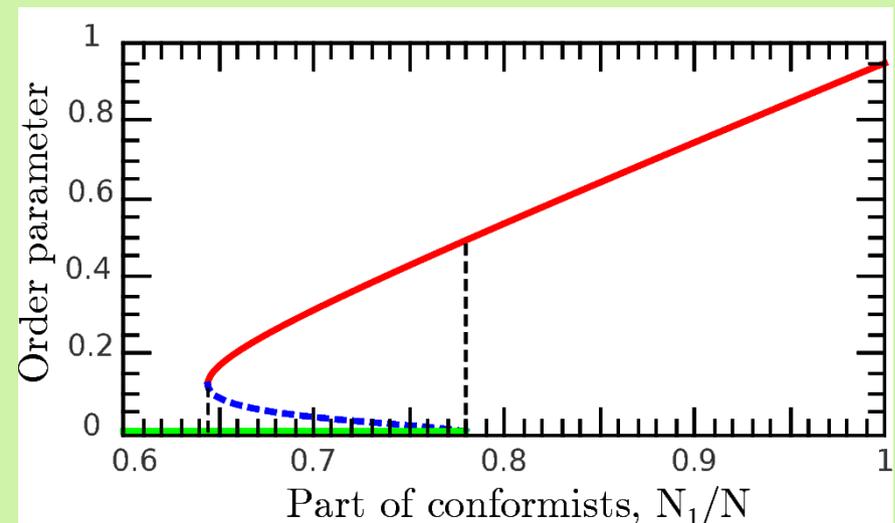
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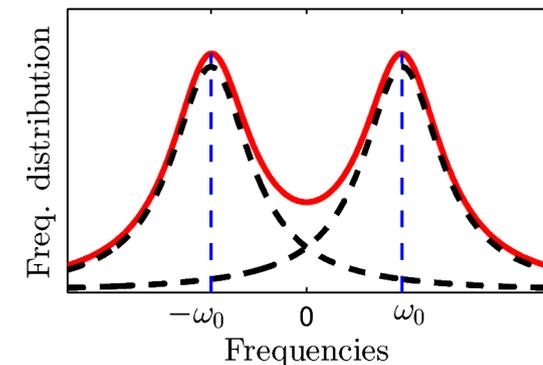
Synchronization bistability

Other examples for phase oscillators:

- Bimodal frequencies distributions¹
- Kuramoto-Sakaguchi model²
- Scale free network³

$$\dot{\theta}_j = \omega_j + \frac{K}{N} \sum_{k=1}^N \sin(\theta_k - \theta_j)$$

Frequencies distributed by the sum of two Lorentz distr.



[1] E. Martens et al., "Exact results for the Kuramoto model with a bimodal frequency distribution", Phys. Rev. E, (2009)

[2] O. E. Omel'chenko and M. Wolfrum, "Nonuniversal transitions to synchrony in the Sakaguchi-Kuramoto model", Phys. Rev. Lett., (2012)

[3] J. G. Gardenes et al., "Explosive Synchronization Transitions in Scale-Free Networks", Phys. Rev. Lett. (2011)

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$$\dot{\theta}_j = \omega_j + \frac{K}{N} \sum_{k=1}^N \sin(\theta_k - \theta_j + \alpha)$$

Some specific unimodal frequencies distributions

[1] E. Martens et al. ,“Exact results for the Kuramoto model with a bimodal frequency distribution”, Phys. Rev. E, (2009)

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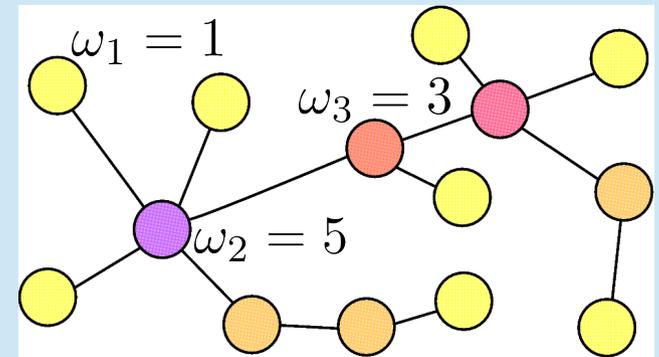
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Probability to be attached to k oscillators $\sim k^{-\gamma}$



$$\dot{\theta}_j = \omega_j + K \sum_{k=1}^N A_{ij} \sin(\theta_k - \theta_j)$$

$$\omega_j \sim \# \text{ attached oscillators}$$

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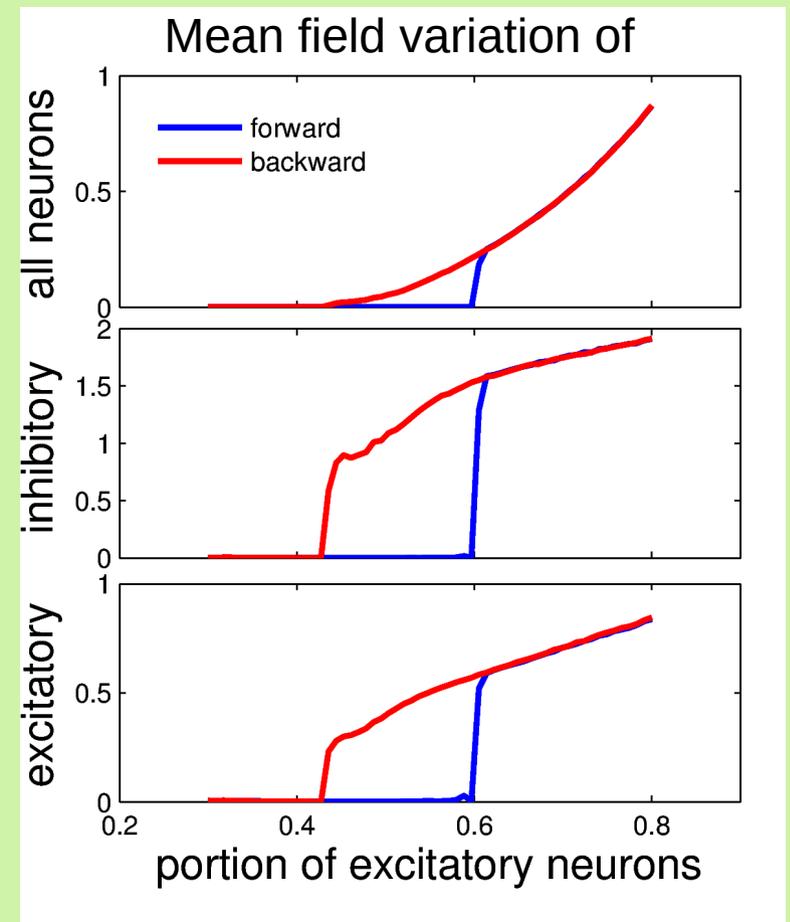
FitzHugh-Nagumo synaptically coupled neurons [“conformists-contrarian” analogue]:

$$\begin{aligned}\dot{v}_j &= v_j(1 - v_j^2/3) - u_j + I + I_j^{(syn)}, \\ \dot{u}_j &= \varepsilon_j(b_0 + b_1 v_j - u_j),\end{aligned}$$

$$I_j^{(syn)} = (v_j^{(0)} - v_j) \frac{K_j}{N} \sum_k \frac{1}{1 + \exp(-(v_k - v_T)/\Delta)}$$

k 'th neuron acts on j 'th only when exceeds some threshold.

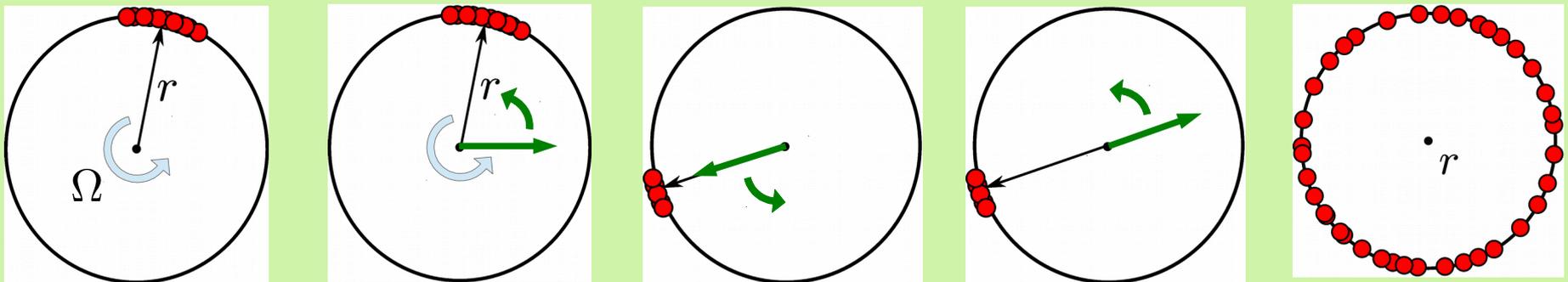
	Excitatory	Inhibitory
K_j	0.425	1.275
$v_j^{(0)}$	2.5	-2.5



Algorithms

- Systems synchronization and phase reversion by π with external periodic force

$$F(t) = a \cos(\Omega t + \phi(t))$$



The moment of external force disconnection should be determined empirically.

- Stimulate system with high frequency periodic signal with decaying amplitude

$$F(t) = a \exp(-t/\tau) \cos(\bar{\Omega}t)$$

Results

“Conformists-contrarians” model [phase change]:

$$\dot{\theta}_j = \omega_j + \frac{1}{N} \sum_{k=1}^N K_k \sin(\theta_k - \theta_j) + a \cos(\Omega t - \theta_j)$$

System:

Coupling $K_k \in \{-3, 1\}$

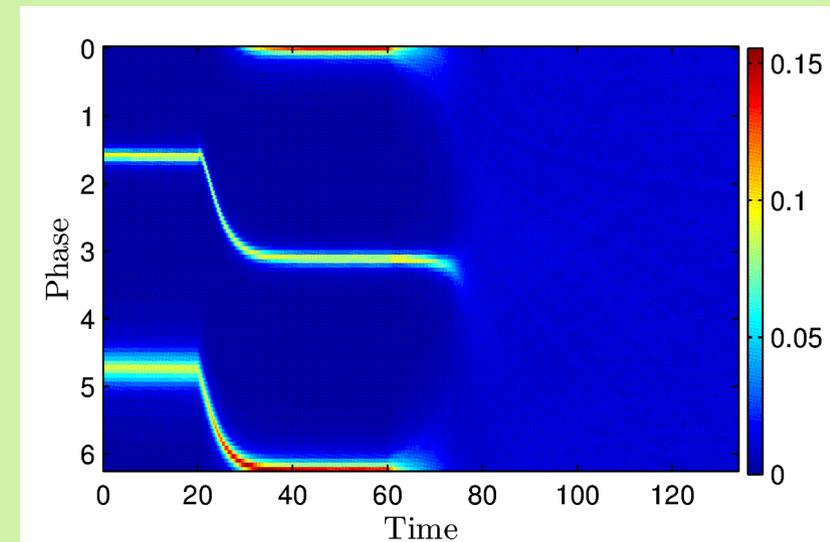
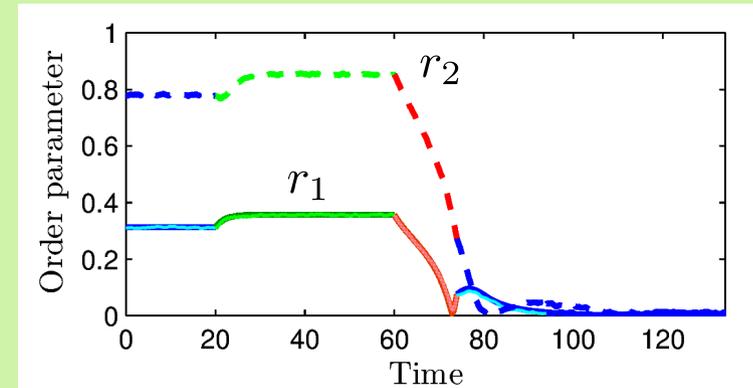
Number of oscillators $N = 10000$

Central frequency $\omega_0 = 0$

Force parameters:

Amplitude $a = 0.2$

Frequency $\Omega = 0$



Results

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$$\dot{\theta}_j = \omega_j + \frac{1}{N} \sum_{k=1}^N K_k \sin(\theta_k - \theta_j) + a \exp(-t/\tau) \cos(\Omega t - \theta_j)$$

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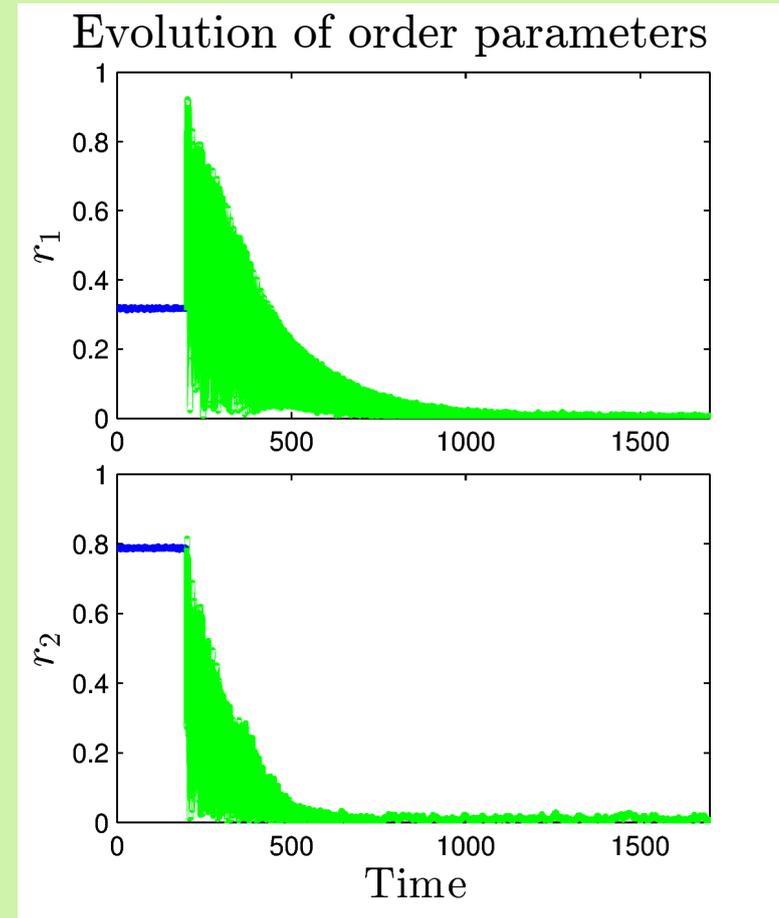
Central frequency $\omega_0 = 0.3$

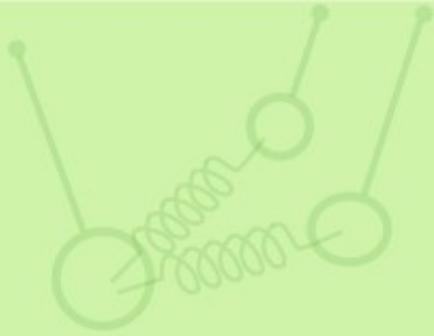
Force parameters:

Amplitude $a = 4$

Frequency $\Omega = 2$

Decay time $\tau = 200$





Results

FitzHugh-Nagumo “conformists-contrarian” model:

Number of neurons: $N = 20000$

Mean field frequency: $\omega_0 \approx 2\pi/75$

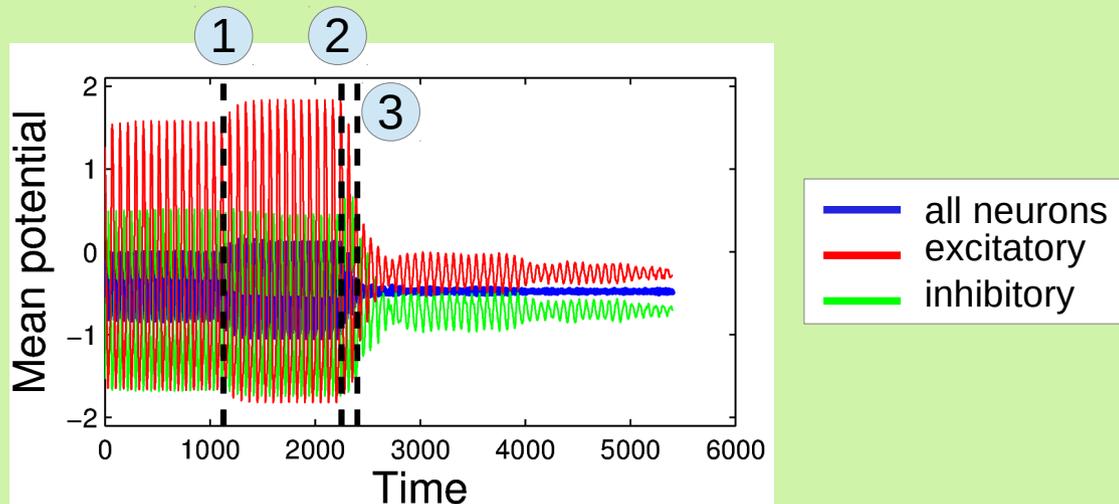
Results

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Phase change: $F(t) = a \cos(\Omega t + \phi(t))$



External force parameters:

$$a = 0.05, \Omega = 2\pi/75$$

Results

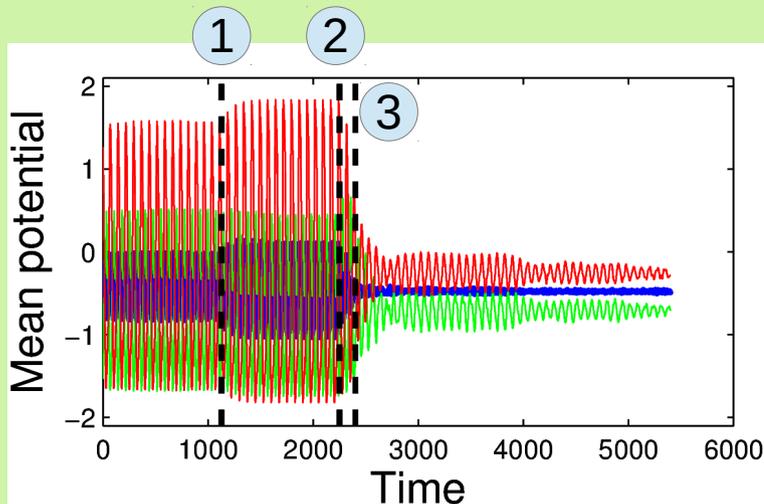
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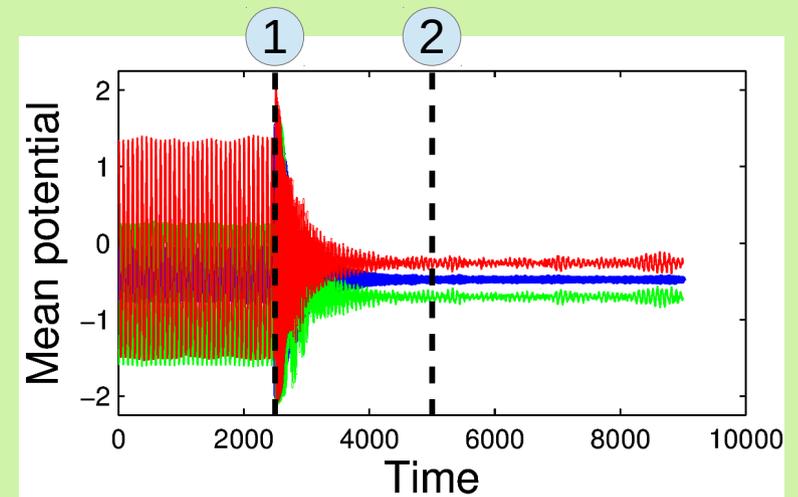
Phase change: $F(t) = a \cos(\Omega t + \phi(t))$

Decaying periodic force: $F(t) = a \exp(-t/\tau) \cos(\Omega t)$



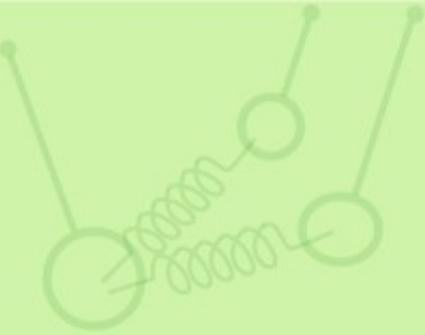
External force parameters:

$$a = 0.05, \Omega = 2\pi/75$$



External force parameters:

$$a = 1.5, \Omega = 2\pi/10, \tau = 500$$

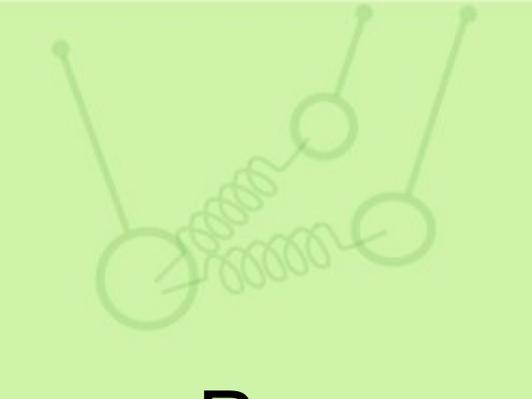


Conclusion

- Proposed algorithms are able to drive synchronized bistable systems to desynchronized state

Further work and unanswered questions:

- How frequent is investigated systems in the nature?
- How to improve algorithms stability? (dependence on number of oscillators, noise, ...)
- How to estimate control parameters?
- etc...



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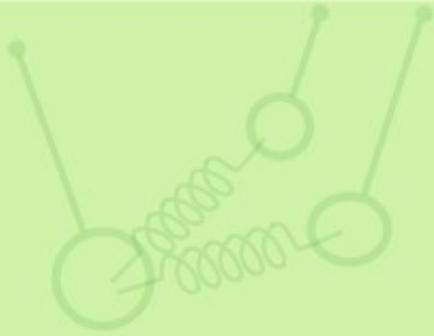
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Thank you for attention!

Acknowledgments

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The end

